

Geometric Transformations in Quantum Field Theory: Implications for the Spectral Analysis of Spacetime

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Abstract—Through the Fukaya conjecture and the wrapped Floer cohomology, the correspondences between paths in a loop space and states of a wrapping space of states in a Hamiltonian space (the ramification of field in this case is the connection to the operator that goes from TM to T^*M) are demonstrated where these last states are corresponding to bosonic extensions of a spectrum of the space-time or direct image of the functor Spec, on space-time. This establishes a distinguished diffeomorphism defined by the mapping from the corresponding loops space to wrapping category of the Floer cohomology complex which furthermore relates in certain proportion D -branes (certain D -modules) with strings. This also gives to place to certain conjecture that establishes equivalences between moduli spaces that can be consigned in a moduli identity taking as spacetime the Hitchin moduli space on G , whose dual can be expressed by a factor of a bosonic moduli spaces.

I. INTRODUCTION

I

N symplectic geometry, the Lagrangian submanifolds are invariants under rescaling of the cotangent fibers or more generally asymptotically invariant which determine diffeomorphisms between the said cotangent classes under homotopies, establishing a cohomology of Floer cohomology type where wrapping states are in correspondence with paths in certain based loop space. By the Fukaya conjecture, this correspondence can be carried to the equivalences context between categories whose objects are Lagrangian submanifolds of the complex Riemannian manifold (complex model of the space-time), and whose morphisms are Floer chain groups of the type $\text{Hom}(L_0, L_1) \cong FC(L_0, L_1)$. Likewise, if $L \subset T_x^*$, has cotangent fiber then the A_∞ structure on

$CW^*(L, L)$, should be quasi-isomorphic to the dg algebra structure on $C_\infty(\square_x)$, where \square_x is the based loop space of (Z, x) , [1]; that is to say, there is a minimal path from (Z, x) , to

□

(L, x) , taking a connection with ramification.

From a point of view of the physics, this could have implications in the process of explanation of the relation between movement and energy, problem planted from Huygens and Lagrange to the dynamical problems, which requires the integration of energy and movement through certain integral operators that can be induced to the micro-

local structure (fine structure) of the Lagrangian submanifolds to the QFT context using the TFT and some tools of cohomology to obtain the equivalences in duality of these Lagrangian objects. From an algebraic point of view, using commutative rings, it can be induced to a scheme that establishes the relations between moduli problems and algebras involved in the reduction problem of field equations to some cases in field theory. Is it feasible to determine symmetry between energy and time? We can answer categorically no! However, realizing a re-evaluating of the question in the framework of the ramified field, could determine an equivalence, not in a purely algebraic context and the continuous mappings only, but in some differential applications that born of several dualities, as have been mentioned (for example, the Langlands duality [2], [3]). Likewise, this happens where their deformable images are of certain applications [8], [9] between differential operators in a holomorphic framework that can be determined establishing actions on said elements (that are ramifications). The actions from loop groups obtained in the construction process of cycles of the space-time [3] are the ramifications.

II. SPECTRUM OF THE COMMUTATIVE RINGS: THE IMAGE OF A

SPIN MANIFOLD

Let X , be a scheme as studied by Milnor [4], more specifically, the scheme given from the derived categories D_X , whose sheaves I , are coherent sheaves of ideals on X , then the transformation that we define is the morphism $\square: X^\sim \rightarrow \square X$, such that $\square_{\square} I_{OX^\sim}$, is an invertible sheaf. Here O_{X^\sim} , is the structure of the sheaf of X^\sim . For this way, morphisms from schemes to affine schemes must be understood in the ring homomorphisms context by the following adjoint contravariant pair (X, A) : For every scheme X , and every commutative ring A , we have the natural equivalence [1], [9]:

$$\text{Hom}_{\text{Schemes}}(X, \text{Spec}(A)) \cong \text{Hom}_{\text{CRing}}(A, O_X(X)), \quad (1)$$

due to that A , is an initial object in the rings category before the functor action Spec, in the schemes category, and their



final object is $\text{Spec}(A)$, which means the spectrum of commutative rings category. The character Spec is the functor “spectrum”.

Likewise, we consider the Axion particle which must be a final element or object in the space $\text{Spec}(A)$. **Definition 1.** A spin manifold in the image of $\text{Spec}(A)$ of the space-time, is an axion.

An axion is not a spin manifold; however, this is the result of a transformation “blow up” in the space-time of a spin manifold in QFT-frame. Likewise, the axion can be determined and defined as an image of a spin manifold under transformation rules in the Universe modeled as a complex Riemannian manifold (of Kahlerian type) whose dimension is $2n$ [7]-[9]. Thus, in moduli problems, we can obtain equivalences in duality of the objects in A , and $\text{Spec}(A)$, and their images in a complex Riemannian manifold.

If we consider as model of the space-time the complex Riemannian manifold $M : \mathbb{C} \setminus_{\mathcal{H}}(G, C)$, the equivalences can be defined in D -branes (that is to say D -modules), and strings level through of the equivalences established in Theorem 7.1 [5], having that, under certain duality, $\mathbb{C} \setminus_{\mathcal{H}}(G, C)$, is composed for objects of a derived category of D -modules on $B_G(C)$, (here $B_G(C)$, as category is a category of Hecke eigensheaves) where this last is a category of the G -equivariant D -modules where C , is a complex whose scheme Y (that are orbifolds of a CY-manifold, which is a Calabi-Yau manifold. This is a spin manifold if establishes equivalences between

D -modules of $B_G(C)$, and D -branes in the moduli space $\mathbb{C} \setminus_{\mathcal{H}}(G, C)$ is the spectrum of G -equivariants D -modules (C , is a complex of certain special sheaf of holomorphic G -bundles (of certain eigen-sheaf of Hecke)). Then can be obtained a spectrum of the space-time.

In the study of the algebraic geometry and the theory of complex analysis on manifolds, the coherent sheaves are a specific class of sheaves having the particularity of establishing properties of differential operators linked to the geometrical properties of the underlying space where these are defined. The geometrical information codification of coherent sheaves is realized with reference to a sheaf of rings.

III. TOPOLOGICAL QUANTUM DIFFEOMORPHISMS IN FIELD THEORY

Using the model in Fig. 1, [1], let the corresponding noncompact Lagrangian submanifolds L , as homotopy of L_0 , where L_0 , is the Lagrangian submanifold before of the field action given by one particle, to know [4]:

$$HW^*(L_0, L_0) \cong H(CW^*(L_0, L_0)), \tag{2}$$

We consider the Fukaya conjecture taking as cotangent fiber $L \subset T_x^*$, to the A_{∞} -structure on complex $CW^*(L_0, L_0)$, which is isomorphic to the algebra dg -structure on the deformed category $C_{\infty}(\mathbb{C}, x)$, of the loop space (Z, x) , with Z , a differentiable manifold which has a normalized geodesic flow accord to their Poincaré section [2], [6].

Remark 1. A free loop space or loops space $\Omega X \subset C(S^1, X)$, of the topological space X , is the space conformed for loops from the unit circle S^1 , to X , endowed with compact-open topology.

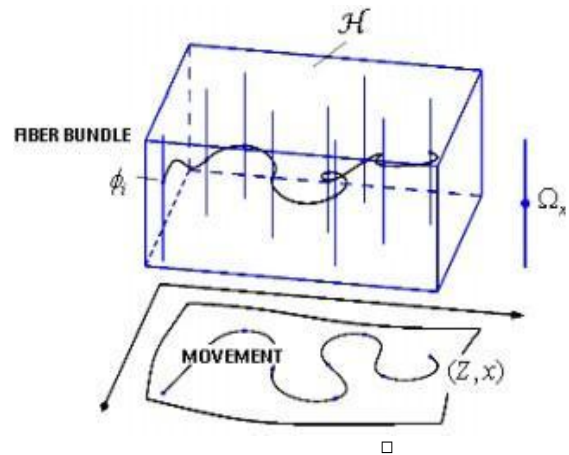


Fig. 1 Correspondences between states (L, x) , of the Hamiltonian space \mathcal{H} , and points (Z, x) , of the loop space

Due to the TFT, the flow is a geodesic flow, which is conformed for trajectories defined by the variation principle. The flow is created from the micro-states that have been measured by gauge fields [1] in a Hamiltonian space as state space [8], [9].

From the conjecture 1, given in [6] the movement is spread ramification from the energy-vacuum mechanism through paths of particle actions given by [7]

$$\square_{\mathcal{O}_c}(\square) \cong \int_M \square_{\mathcal{O}_c}(\square(x)) d\square(x), \tag{3}$$

The Bulnes’s operator \mathcal{O}_c [7], involves a connection

$\square_t(X_p), \square_p \square M, X_p \square T_p(M)$ of the tangent bundle of the space of trajectories $\square(\square) \square \mathbb{R}^3 \square I_t$, such that the map

$$X \square \square_x(1), \tag{4}$$

Indeed, is the affine connection describes as:

types [1]. Thus, considering the functoriality scheme of rings established in (1) and having the connection scheme to the space-time \mathbb{M} , ramified field [1], [9] we can enounce the following generalizing of (11) given by the Theorem.

Theorem 2. (F. Bulnes) [1], [12], [13]. If we consider the category $M_{K^r}(g^{\wedge}, Y)$, then a scheme of their spectrum $W(H)$, where Y , is a Calabi-Yau manifold comes given as:

$$\text{Hom}_{g^{\wedge}}(X, V_{\text{criticalDef}}) \square \text{Hom}_{\text{Loc } L_G}(V_{\text{critical}}, M_{K^r}(g^{\wedge}, Y)), \quad (11)$$

Proof. [1], [2], [14].

The spectrum in (11) is a Verma module of critical level. In the case to consider a homotopy category $H \circ (C)$, of a differential graded category C , then the isomorphism given in (11) has in some time, the same objects as C , but their morphisms are defined by the identity:

$$\text{Hom}_{H \circ (C)}(L_o, L) \square H \circ \text{Hom}_C(L_o, L), \quad (12)$$

which is doing re-counting to image of spectrum over all space-time \mathbb{M} , considering this as the Hitchin space [15] $\backslash \text{ }_H(G, C)$, to different stratus of dimensional spaces, we can derive an identity of the moduli spaces considering their different geometrical and physical stacks [16].

We demonstrate the commutative scheme given in (10), using the fact that

$$\text{emb}_{\mathbb{R}^{\square 1}}(O_c(\square))(\square) \square \text{Diff}(\square)(x(t)), \quad (13)$$

Using arguments of homotopy [17], and the natural chain mappings or diffeomorphism of the type (in homological algebra, the homotopy category of chain complexes in an additive category C , is a framework for working with chain homotopies and homotopy equivalences [17]. This is an intermediate between the category of chain complexes of C , and the derived category of C , when C , is Abelian; unlike the former it is a triangulated category, and unlike the latter its formation does not require that C , is Abelian):

$$\square : C(L_o, L_i) \square WH(L_o, L_i) \square WH(L), \quad (14)$$

where $WH(L) \square W(H)$, we obtain images $C \square WH$, which complete the diagram (10).

In the general case, that is to say, in all space \mathbb{M} , will be necessary use the natural chain mapping

$$\square : CW(F) \square C_{\square^*}(\square(M)), \quad (15)$$

here $F \square T_q^*M \square T^*M$.

V. CONCLUSION

The spectrum of the deformed derived category given for can be viewed as an aspect of one version of the Floer cohomology type and their Lagrangians (Lagrangian submanifolds) that are the object (points of the supermanifold defined by the Hitchin moduli space $\backslash \text{ }_H(G, C)$) in the other class of the homomorphism $L_o \square L_i$, that includes the ring structure [2] to preserve commutativity in the diagram (10). This is the principal conclusion derived of the Fukaya conjecture, and schematized through (10), which is commutative diagram to geometrical Langlands ramification of the operator $O_c : TM \square T^*M$. Likewise, this ramification is given by a connection deduced from the mentioned operator whose states are in the modules of the category $WH(L) \square W(H)$, where each Lagrangian submanifold is a kernel to the Floer homology complex, that is to say, $WH(L) \square 0$. Then, a more general sense, that is to say, to category of Lagrangians $F \square T_q^*M \square T^*M$, we have (15).

Then in wrapped Floer cohomology we have a field equation as for example $W \square \square 0$, (Weyl equation) or the Dirac equation

$$D \square \square 0.$$

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